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INTERACTIONS OF SECONDARY PARTICLES
WITH THORIUM SAMPLES IN THE SETUP **QUINTA**
IRRADIATED WITH 6-GeV DEUTERONS

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Взаимодействие вторичных частиц с образцами тория
на установке «Квинта» при облучении 6-ГэВ дейтронами

Сборка из природного урана (512 кг) «Квинта» облучалась 6-ГэВ дейтронами. Образцы ^{232}Th были расположены на центральной оси установки «Квинта». При анализе спектров гамма-лучей, испущенных облученными образцами ^{232}Th , было идентифицировано более ста радионуклидов продуктов реакций, в том числе и ^{233}U . Определено отношение массы образующегося ^{233}U к массе ^{232}Th . Результаты эксперимента сравниваются с модельными расчетами по методу Монте-Карло (код FLUKA).

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Interactions of Secondary Particles with Thorium Samples
in the Setup QUINTA Irradiated with 6-GeV Deuterons

The natural uranium assembly, QUINTA, was irradiated with 6-GeV deuterons. The ^{232}Th samples were placed at the central axis of the setup QUINTA. The spectra of gamma rays emitted by the activated ^{232}Th samples have been analysed, and more than one hundred nuclei produced have been identified. For each of those products, reaction rates have been determined. The ratio of the weight of produced ^{233}U to that of ^{232}Th is presented. Experimental results were compared with the results of Monte Carlo simulations by the FLUKA code.

The investigation has been performed at the Dzhelepov Laboratory of Nuclear Problems, JINR.

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1. INTRODUCTION

Accelerator Driven Systems (ADSs) are suggested as a means for safe, cost-effective energy production and nuclear waste transmutation [1–4]. An ADS consists of a high-current, high-energy accelerator coupled with a subcritical nuclear assembly. During the past several years, such studies have been conducted with accelerated particle beams at the Nuclotron of the Joint Institute for Nuclear Research (JINR) in the framework of the international collaboration “Energy plus Transmutation of Radioactive Waste”. This program has carried out a large number of experiments with the subcritical uranium target QUINTA [5–12], as well as the lead-graphite target GAMMA-3 [13, 14]. Several experiments were conducted using a solid lead target GENERATOR [15–17] at the proton beam of the JINR Phasotron accelerator. ^{232}Th is a fissionable material by fast neutrons, but neutron capture reaction is also important in order to study the breeding efficiency of fissile ^{233}U in a thermal-spectrum ADS. In this paper we present new experimental data and a comparison with the calculations by the FLUKA [18, 19] code. For thorium samples, located at the central axis of the setup QUINTA, the reaction rates for generated radionuclides were measured and are presented in this paper.

2. STRUCTURE OF THE SETUP QUINTA

Uranium assembly QUINTA is presented in Fig. 1. It consists of five sections of hexagonal shaped aluminum containers with an inscribed circular hole of diameter 28.4 cm. The containers are filled with cylindrical rods of natural uranium, having a sealed aluminum shell (external dimensions: 3.6 cm diameter, 10.4 cm in length, and 1.72 kg uranium mass). The containers are made of 5-mm-thick aluminum. The first section, facing the deuteron beam, contains 54 uranium rods and has a central beam window, 80 mm in diameter, installed in order to reduce its albedo and reduce the leakage of neutrons from the target. Four subsequent sections are structurally identical and contain 61 uranium rods. The mass of the natural uranium in each of these sections is 104.92 kg, and the

Fig. 1. General view of the QUINTA setup

total mass of uranium in the entire target is 512.56 kg. The filling factor of the 2nd, 3rd, 4th, and 5th uranium sections is about 0.8, and is ~ 0.6 of the entire assembly.

3. EXPERIMENT

During the experiment, the ^{232}Th samples were conventionally marked as 9Th, 10Th, 11Th and 12Th, masses of the samples are respectively 0.1236, 0.1242, 0.1355 and 0.1402 g. The 9Th sample was placed on the central axis of the setup QUINTA, between the first and second sections at a distance of 12.1 cm, the 10Th sample between the third and fourth sections at a distance of 25.2 cm, the 11Th sample between the third and fourth sections at a distance of 38.3 cm, and 12Th sample between the fourth and fifth sections at a distance of 51.4 cm from the front end. The total number of incident deuterons colliding with the target was determined by the standard method of activation of aluminum foil by the reaction $^{27}\text{Al}(d, x)^{24}\text{Na}$. After 975 minutes irradiation the total number of deuterons incident on the target was $1.93(2)\text{E} + 13$. The technique used to determine the total number of incident deuterons on target is described in detail in Ref. [5]. The γ -ray spectra of the samples were measured with three

HPGe detectors manufactured by ORTEC (two detectors with relative efficiencies of 28% and 33% and the energy resolutions of 1.9 and 1.8 keV, respectively, at the 1.33-MeV ^{60}Co line) and CANBERRA (one detector with a relative efficiency of 19%, and an energy resolution of 1.8 keV at the 1.33-MeV ^{60}Co line). For each sample, from 8 to 13 γ -ray spectra were measured with different time intervals and the cooling time of the first spectrum ranged from 120 to 160 min. Energy and efficiency calibrations of the detectors were performed using a set of the γ -ray standards (^{54}Mn , ^{57}Co , ^{60}Co , ^{88}Y , ^{113}Sn , ^{133}Ba , ^{137}Cs , ^{139}Ce , ^{152}Eu , ^{228}Th , ^{241}Am).

The primary analysis of the measured γ -ray spectra was performed using the DEIMOS32 code [20]. The program allows determining the areas under the peaks and their positions (channel number). After that, using a software package [21], the spectra were calibrated for energy, corrected for the detector efficiency, and separate γ -ray lines of the product nuclei were identified as formed in the samples as a result of interactions with secondary neutrons, protons or deuterons. Experimental count rates of the individual γ -ray transitions were corrected for the nuclear decay during the irradiation as well as for the self-absorption for the γ -rays measured, for the geometric dimensions of the samples, for the true coincidence summing [22], for the beam interruptions during irradiation and the variations in intensity of the deuteron beam (based on the on-line measurements with fast ionization chambers). All these procedures are described in detail in Refs. [21, 23, 24].

For determining the experimental reaction rates, the following equation was used [24]:

$$R(A_r, Z_r) = \frac{Q(A_r, Z_r)}{N_t N_d}, \quad (1)$$

where $Q(A_r, Z_r)$ is the production rate of the radioactive nucleus (A_r, Z_r) , N_t the number of atoms in the sample, and N_d the number of incident deuterons on the target.

4. MONTE CARLO SIMULATIONS

The FLUKA hadron–nucleon interaction models are based on resonance production and decay at energies below a few GeV, and on the Dual Parton model at energies above. Two models are also used in hadron–nucleus interactions. At momenta below 3–5 GeV/c, the PEANUT package [25, 26] includes a very detailed Generalised Intra-Nuclear Cascade (GINC) and a preequilibrium stage, while at high energies the Gribov–Glauber multiple collision mechanism is included in a less refined GINC. Both modules are followed by equilibrium processes: evaporation, fission, Fermi break-up, γ -ray de-excitation [27, 28]. Inelastic cross sections for hadron–hadron interactions are represented by parameterised fits

Fig. 2. Total fluence of secondary neutrons (a), total fluence of secondary protons (b) and total fluence of primary and secondary deuterons in the setup QUINTA (c). Shown is side view of the setup geometry cut from the center by Z axis

based on available experimental data [41]. For hadron–nucleus interactions, a mixture of tabulated data and parameterised fits is used [29–33]. Elastic and charge-exchange reactions are described by phase-shift analyses and eikonal approximation. Ion-induced nuclear interactions are described for energy between 0.1 and 5 GeV/nucleon with modified Relativistic Quantum Molecular Dynamics (RQMD) model [34–36] and for energy below 0.1 GeV/nucleon with Boltzmann Master Equation (BME) theory [37–39].

Figure 2 shows (a) the total fluence of secondary neutrons, (b) the total fluence of secondary protons and (c) the total fluence of primary and secondary deuterons in the QUINTA setup (cut geometry from the center by Z axis is shown conventionally). By Y and Z axes the size of the setup is shown. Fluence was produced by one 6-GeV deuteron and averaged by the X axis. A calculated double differential fluence of secondary neutrons for the position of the thorium samples is presented in Fig. 3. As can be seen from the picture, the biggest fluence of neutrons is for sample 10Th and the smallest fluence is for samples 9Th and 12Th at energies below 10 MeV and at energies above 10 MeV smallest fluence is only for sample 9Th. The deuteron beam parameters used in the simulations were determined experimentally in Ref. [5]. Coordinates of the center of the beam are $X_C = 2.0$ cm, $Y_C = -0.1$ cm and FWHM-distribution is $\text{FWHM}_X = 3.9$ cm, $\text{FWHM}_Y = 3.1$ cm. The dependence of reaction rate on mass of the residual isotopes for the samples 9Th, 10Th, 11Th and 12Th is shown in Fig. 4, in units of $\text{nuclei/cm}^3 \cdot \text{deuteron}$. The products of evaporation, fragmentation, fission (with low and high energy particles), spallation and quasi-elastic reactions in the thorium samples can be seen in Fig. 4.

Fig. 3. Double differential fluence of secondary neutrons for the position of the samples 9Th, 10Th, 11Th and 12Th

Fig. 4. Dependence of reaction rate on mass of the residual isotopes for the samples 9Th, 10Th, 11Th and 12Th

In the results of simulations, detailed information is taken about the number of fissions in each of the natural uranium bars of the QUINTA setup. In Table 1 the number of fissions for natural uranium bars per one 6-GeV deuteron is given and bars are marked as shown in Fig. 5. The total number of fissions in the setup generated by one 6-GeV deuteron is 51.3, and 14.9 fissions of them are generated with high-energy ($E > 20$ MeV) particles. From the fissions with high-energy ($E > 20$ MeV) particles, 0.34 fissions are generated by deuterons, 1.7 by protons, 11.3 by neutrons, 0.37 by positive pions, and 1.22 by negative pions. The total number of fissions in the 1st section is 2.7, in the 2nd section is 16.4, in the 3rd section is 16.0, in the 4th section is 10.7, and in the 5th section is 5.6. The number of fissions with high-energy ($E > 20$ MeV) particles in the 1st section is 0.52, in the 2nd section is 4.96, in the 3rd section is 4.66, in the 4th section is 3.08, and in the 5th section is 1.72.

Table 1. Number of fissions for natural uranium bars per one 6-GeV deuteron

Number of bars	1st section	2nd section	3rd section	4th section	5th section
U01	1.94E-02	4.43E-02	5.63E-02	4.38E-02	2.49E-02
U02	2.56E-02	6.40E-02	8.09E-02	6.19E-02	3.50E-02
U03	3.06E-02	8.21E-02	1.03E-01	7.88E-02	4.46E-02
U04	3.18E-02	8.73E-02	1.12E-01	8.68E-02	4.98E-02
U05	2.75E-02	7.46E-02	9.84E-02	7.84E-02	4.63E-02

Table 1 continued

Number of bars	1st section	2nd section	3rd section	4th section	5th section
U06	2.30E-02	5.26E-02	6.54E-02	4.96E-02	2.79E-02
U07	3.42E-02	8.81E-02	1.07E-01	7.92E-02	4.38E-02
U08	4.39E-02	1.30E-01	1.54E-01	1.12E-01	6.15E-02
U09	5.09E-02	1.65E-01	1.97E-01	1.43E-01	7.79E-02
U10	4.85E-02	1.58E-01	1.97E-01	1.48E-01	8.44E-02
U11	3.72E-02	1.12E-01	1.47E-01	1.16E-01	6.90E-02
U12	2.33E-02	5.30E-02	6.56E-02	4.94E-02	2.76E-02
U13	4.01E-02	9.87E-02	1.16E-01	8.42E-02	4.59E-02
U14	6.23E-02	1.76E-01	1.96E-01	1.35E-01	7.12E-02
U15	7.79E-02	2.90E-01	3.06E-01	2.02E-01	1.05E-01
U16	8.35E-02	3.54E-01	3.86E-01	2.63E-01	1.39E-01
U17	6.82E-02	2.57E-01	3.13E-01	2.34E-01	1.33E-01
U18	4.49E-02	1.45E-01	1.91E-01	1.52E-01	9.09E-02
U19	2.08E-02	4.58E-02	5.65E-02	4.29E-02	2.41E-02
U20	3.66E-02	9.08E-02	1.06E-01	7.68E-02	4.16E-02
U21	6.47E-02	1.82E-01	1.95E-01	1.31E-01	6.86E-02
U22		4.22E-01	3.79E-01	2.32E-01	1.14E-01
U23		1.02E+00	7.84E-01	4.33E-01	1.99E-01
U24	1.79E-01	8.10E-01	8.28E-01	5.51E-01	2.83E-01
U25	8.23E-02	3.30E-01	4.06E-01	3.14E-01	1.87E-01
U26	4.42E-02	1.43E-01	1.91E-01	1.56E-01	9.64E-02
U27	1.59E-02	3.34E-02	4.17E-02	3.25E-02	1.87E-02
U28	2.96E-02	6.82E-02	8.23E-02	6.10E-02	3.37E-02
U29	5.54E-02	1.36E-01	1.52E-01	1.05E-01	5.59E-02
U30		3.19E-01	2.97E-01	1.86E-01	9.31E-02
U31		1.29E+00	7.75E-01	3.87E-01	1.72E-01
U32		2.71E+00	1.81E+00	8.96E-01	3.68E-01
U33	1.74E-01	7.95E-01	9.03E-01	6.56E-01	3.62E-01
U34	6.66E-02	2.46E-01	3.19E-01	2.56E-01	1.58E-01
U35	3.36E-02	1.02E-01	1.40E-01	1.18E-01	7.37E-02
U36	2.04E-02	4.52E-02	5.58E-02	4.26E-02	2.38E-02
U37	3.63E-02	8.92E-02	1.05E-01	7.62E-02	4.14E-02
U38	6.39E-02	1.79E-01	1.91E-01	1.30E-01	6.73E-02
U39		4.06E-01	3.70E-01	2.27E-01	1.12E-01
U40		9.38E-01	7.39E-01	4.16E-01	1.91E-01
U41	1.61E-01	7.59E-01	7.80E-01	5.21E-01	2.67E-01
U42	7.98E-02	3.22E-01	3.97E-01	3.06E-01	1.82E-01

Table 1 continued

Number of bars	1st section	2nd section	3rd section	4th section	5th section
U43	4.31E-02	1.40E-01	1.88E-01	1.54E-01	9.51E-02
U44	2.29E-02	5.16E-02	6.37E-02	4.81E-02	2.70E-02
U45	3.93E-02	9.68E-02	1.13E-01	8.25E-02	4.46E-02
U46	6.11E-02	1.71E-01	1.90E-01	1.32E-01	6.95E-02
U47	7.62E-02	2.76E-01	2.94E-01	1.96E-01	1.01E-01
U48	8.05E-02	3.35E-01	3.68E-01	2.52E-01	1.34E-01
U49	6.51E-02	2.46E-01	3.02E-01	2.26E-01	1.29E-01
U50	4.32E-02	1.40E-01	1.85E-01	1.48E-01	8.87E-02
U51	2.20E-02	5.02E-02	6.29E-02	4.79E-02	2.70E-02
U52	3.30E-02	8.46E-02	1.03E-01	7.66E-02	4.25E-02
U53	4.23E-02	1.24E-01	1.48E-01	1.08E-01	5.89E-02
U54	4.88E-02	1.57E-01	1.88E-01	1.37E-01	7.46E-02
U55	4.63E-02	1.50E-01	1.88E-01	1.42E-01	8.08E-02
U56	3.53E-02	1.07E-01	1.40E-01	1.12E-01	6.61E-02
U57	1.80E-02	4.11E-02	5.21E-02	4.07E-02	2.33E-02
U58	2.39E-02	5.99E-02	7.53E-02	5.80E-02	3.32E-02
U59	2.86E-02	7.63E-02	9.68E-02	7.38E-02	4.18E-02
U60	2.94E-02	8.17E-02	1.05E-01	8.15E-02	4.65E-02
U61	2.56E-02	6.94E-02	9.23E-02	7.38E-02	4.34E-02

Fig. 5. Marking of ^{nat}U bars of the QUINTA setup for the definition of fission numbers

5. EXPERIMENTAL RESULTS AND COMPARISONS WITH MONTE CARLO SIMULATIONS

Table 2 gives the cumulative reaction rates for residual radionuclides obtained from the experiment. As can be seen from the table, for the sample 10Th the biggest values of reaction rates are for almost all residual nuclei and the smallest values of reaction rates are for sample 12Th. ^{233}Pa is produced in the reaction $^{232}\text{Th}(n, \gamma)^{233}\text{Th} (\beta^- \text{ decay, } T_{1/2} = 22.3 \text{ min}) \rightarrow ^{233}\text{Pa} (\beta^- \text{ decay, } T_{1/2} = 26.967 \text{ day}) \rightarrow ^{233}\text{U}$. Ratio of the weight of produced ^{233}U to ^{232}Th for the sample 9Th is $1.39(4)\text{E-}12$, for 10Th is $3.51(10)\text{E-}12$, for 11Th is $2.39(8)\text{E-}12$, and for 12Th equals $1.48(4)\text{E-}12$.

Table 2. Experimentally obtained cumulative reaction rates for residual radionuclides

Residual nuclei	Reaction rate ($\text{atom}^{-1} \cdot \text{deuteron}^{-1}$)			
	9Th	10Th	11Th	12Th
Be-7		9.0(11)E-27		
Na-22		7.5(16)E-27		
Na-24	1.47(5)E-27	1.00(2)E-27	3.78(22)E-28	1.25(18)E-28
Mg-28	6.44(48)E-28			
S-38	2.73(61)E-28			
Ar-41			5.00(71)E-28	
K-42	6.11(95)E-28			
K-43	2.09(9)E-27	4.01(5)E-28	2.84(25)E-28	1.08(4)E-27
Sc-43		7.8(14)E-28		
Ca-47		2.07(31)E-28		
Sc-47		5.51(52)E-28		
Sc-48	3.71(25)E-28			
V-48		1.05(18)E-27		
Mn-52		1.90(49)E-27		
Mn-56	1.33(25)E-27			
Ni-56		7.8(36)E-28		
Fe-59		7.5(11)E-28		
Ga-66	9.7(27)E-28	1.22(22)E-27		3.48(43)E-28
Ge-66				1.57(28)E-27
Zn-69m	3.07(18)E-28	3.03(16)E-28	2.22(15)E-28	7.9(11)E-29
As-70		1.50(27)E-27	1.65(29)E-27	
Zn-71m		2.39(24)E-28		
As-71	1.97(19)E-28			
As-72		1.84(70)E-27	6.2(11)E-28	

Table 2 continued

Residual nuclei	Reaction rate (atom ⁻¹ · deuteron ⁻¹)			
	9Th	10Th	11Th	12Th
Ga-72		8.8(19)E-28	1.02(13)E-28	
Ga-73		1.12(19)E-27		
Se-73		1.87(29)E-28		
As-74	7.3(10)E-28	2.10(36)E-27		
As-76		9.77(68)E-28		
Br-76		6.17(96)E-28	2.97(27)E-28	
Ge-77		1.01(14)E-27	4.46(62)E-28	1.43(33)E-28
As-78	7.8(13)E-28			
Rb-81		3.91(29)E-28		
Br-82		9.61(34)E-28	6.78(65)E-28	
Rb-82m		6.87(63)E-28		
Rb-83	2.95(35)E-27	1.47(30)E-27		
Sr-85	1.45(12)E-27			
Kr-85m	2.09(8)E-27	3.92(61)E-27	3.01(28)E-27	1.30(9)E-27
Y-85m		6.1(19)E-28		
Rb-86		1.12(37)E-27		
Y-86		4.08(42)E-28		
Zr-86	4.80(28)E-28	4.48(44)E-28	2.42(16)E-28	
Kr-87			3.08(32)E-27	1.29(12)E-27
Y-87			4.66(89)E-28	
Y-87m	1.22(9)E-27	1.07(5)E-27	5.43(63)E-28	2.55(20)E-28
Sr-87m		1.28(38)E-28		
Kr-88	2.64(57)E-27	4.60(18)E-27	3.04(29)E-27	1.67(17)E-27
Y-88		1.45(18)E-27		
Nb-89m		4.22(88)E-28		
Nb-90	4.02(34)E-28	1.12(3)E-27		
Y-90m	1.06(10)E-27		6.44(65)E-28	2.46(17)E-28
Y-91m		1.29(5)E-26		
Sr-91	2.96(17)E-27	5.90(27)E-27	4.42(17)E-27	2.06(11)E-27
Sr-92	2.86(18)E-27	5.17(16)E-27	3.61(21)E-27	1.80(13)E-27
Y-92	4.66(38)E-27	9.66(70)E-27	6.45(61)E-27	3.21(82)E-27
Y-93	4.95(66)E-27	5.76(40)E-27	5.41(31)E-27	2.08(18)E-27
Mo-93m	3.46(58)E-28	2.72(81)E-28		
Tc-94	3.51(46)E-28			
Tc-94m		1.28(19)E-27		
Ru-94		2.21(30)E-27		

Table 2 continued

Residual nuclei	Reaction rate (atom ⁻¹ · deuteron ⁻¹)			
	9Th	10Th	11Th	12Th
Zr-95		7.47(61)E-27	4.90(32)E-27	
Nb-95		2.87(51)E-27		
Tc-95		5.1(13)E-28		
Ru-95		3.66(27)E-27		
Nb-96	7.96(96)E-28	1.39(14)E-27	8.37(55)E-28	3.96(40)E-28
Tc-96		5.7(10)E-28		
Zr-97	2.36(11)E-27	5.09(16)E-27	3.41(14)E-27	1.69(6)E-27
Nb-97		4.94(17)E-27		
Nb-98m		1.71(22)E-27		
Tc-99m		1.21(5)E-27		
Mo-99	3.48(37)E-27		4.31(30)E-27	2.04(12)E-27
Rh-99		2.83(15)E-27		
Rh-99m		4.5(11)E-28		
Rh-100	4.54(27)E-28	4.62(57)E-28		
Rh-101m	9.87(56)E-28	7.9(12)E-28		
Pd-101		1.40(16)E-27		
Ru-103	3.12(14)E-27	7.35(42)E-27	4.11(17)E-27	1.94(12)E-27
Ru-105	2.95(13)E-27	5.16(11)E-27	3.65(14)E-27	1.57(13)E-27
Rh-105	2.84(21)E-27	5.89(23)E-27	3.82(28)E-27	1.80(12)E-27
In-108		9.0(24)E-28		
Ag-110m		6.2(33)E-27	2.65(28)E-27	
In-110			1.70(16)E-27	
Pd-111m		1.42(6)E-27	9.2(15)E-28	
Ag-111		5.76(72)E-27		
Cd-111m		3.91(28)E-27		
In-111	8.3(12)E-28	6.23(22)E-28	2.67(22)E-28	1.48(43)E-28
Ag-112		6.3(19)E-27		
Ag-113		6.28(71)E-27	4.93(54)E-27	2.26(23)E-27
Sn-113		1.46(34)E-27		
Cd-115	1.55(7)E-27		2.44(16)E-27	1.17(7)E-27
In-115m		1.24(39)E-27		
In-116m		3.94(54)E-27		
Sb-116		6.76(86)E-28		
Sb-116m			1.20(15)E-27	
Te-116		2.15(21)E-27		
Cd-117		9.72(94)E-28	9.21(85)E-28	3.94(52)E-28

Table 2 continued

Residual nuclei	Reaction rate (atom ⁻¹ · deuteron ⁻¹)			
	9Th	10Th	11Th	12Th
Cd-117m		1.56(11)E-27	1.23(15)E-27	
Sb-117	3.65(32)E-27	6.6(11)E-27	3.42(21)E-27	
Sn-117m	1.03(18)E-27			
Sb-118m	3.17(30)E-27	4.50(99)E-28	2.19(25)E-28	1.47(24)E-28
Te-119	4.42(61)E-28	3.00(24)E-28	2.11(20)E-28	
Te-119m		3.04(38)E-28		
Sb-120m		6.13(69)E-28	4.49(62)E-28	1.78(30)E-28
I-120m		2.26(72)E-27		
Te-121		7.04(90)E-28		
Te-121m		3.03(36)E-27		
I-121		7.44(67)E-28		
Sb-122	4.63(70)E-28	7.88(75)E-28	5.01(93)E-28	
Xe-122		2.40(25)E-27		
Te-123m		2.43(91)E-27		
I-123		2.80(28)E-27		
Xe-123		1.09(13)E-27		
Sb-124		1.71(32)E-27		8.7(22)E-28
I-124		7.8(20)E-28	3.99(40)E-28	
Xe-125	1.52(9)E-27	6.78(32)E-28	3.82(23)E-28	1.58(21)E-28
Sb-126		9.16(44)E-28		
I-126		2.17(22)E-27		
Sb-127		1.16(10)E-27		
Xe-127	9.52(91)E-28	1.16(6)E-27	8.2(16)E-28	
Cs-127		1.50(13)E-27	7.6(13)E-27	
Sb-128	3.27(61)E-28	5.65(53)E-28	5.22(55)E-28	3.03(36)E-28
Ba-128		6.3(30)E-28		
Cs-129	2.28(11)E-27			
Sb-129	4.64(65)E-27	1.15(7)E-27	6.95(69)E-28	
Ba-129		1.35(36)E-27		
I-130	2.60(29)E-28		5.16(95)E-28	2.67(28)E-28
I-131		2.56(22)E-27	2.23(17)E-27	9.92(60)E-28
Te-132	1.42(17)E-27	2.32(11)E-27	1.90(8)E-27	9.13(90)E-28
I-132	8.1(12)E-28	2.23(61)E-27	1.37(17)E-27	
Cs-132	1.75(12)E-28	2.85(45)E-27	2.46(14)E-27	1.43(23)E-27
La-132		6.39(99)E-28	4.10(58)E-28	1.19(38)E-28
Ce-132		5.76(44)E-28	2.95(56)E-28	

Table 2 continued

Residual nuclei	Reaction rate (atom ⁻¹ · deuteron ⁻¹)			
	9Th	10Th	11Th	12Th
I-133	1.70(7)E-27	3.16(28)E-27	2.48(9)E-27	1.33(5)E-27
La-133		1.12(23)E-26		
I-134		1.02(15)E-26	6.2(10)E-27	3.96(25)E-27
I-135	1.84(17)E-27	3.39(14)E-27	2.24(10)E-27	1.17(9)E-27
Xe-135	2.29(45)E-27	3.77(43)E-27	3.29(69)E-27	1.64(35)E-27
Cs-135m		1.54(21)E-27		
La-135		8.8(19)E-27		
Ce-135	1.45(13)E-27	4.70(51)E-28	1.30(19)E-27	6.38(93)E-28
Cs-136		4.59(24)E-28		
Cs-138		6.6(10)E-27		
Pr-138m		2.77(69)E-28		
Ba-139				2.24(36)E-27
Ba-140	3.44(58)E-27	6.33(48)E-27	3.65(48)E-27	2.98(60)E-27
La-140	3.92(69)E-28	1.15(30)E-27	5.25(80)E-28	3.91(77)E-28
Ce-141		6.07(50)E-27	6.30(40)E-27	
La-142	2.26(21)E-27		2.92(23)E-27	1.19(11)E-28
Ce-143	2.10(9)E-27	3.50(6)E-27	3.07(11)E-27	1.57(7)E-27
Eu-145		6.7(16)E-28		
Eu-146		1.02(12)E-27		
Eu-147		3.41(48)E-27		
Gd-149	1.52(47)E-27	1.14(16)E-27		
Tb-152		6.61(50)E-28	2.76(48)E-28	
Dy-152	7.37(65)E-28	3.98(31)E-28	2.38(52)E-28	
Dy-155	8.90(44)E-28	6.66(22)E-28	3.60(37)E-28	1.43(24)E-28
Dy-157		8.35(23)E-28	4.61(35)E-28	1.86(27)E-28
Eu-157	9.3(33)E-28			
Er-161	1.44(17)E-27	1.11(7)E-27		3.28(62)E-28
Tm-167		1.08(21)E-27		
Tm-168		1.34(24)E-27	2.38(47)E-27	
Lu-170		2.52(42)E-27		
Hf-170	1.08(8)E-28	1.18(21)E-27	5.05(67)E-28	
Hf-173		1.09(8)E-27	6.60(66)E-28	2.68(44)E-28
Ta-173		2.89(29)E-27		
Ta-174		1.89(20)E-27		
Ta-175			4.43(34)E-27	
Yb-175		1.99(55)E-27		

Table 2 continued

Residual nuclei	Reaction rate (atom ⁻¹ · deuteron ⁻¹)			
	9Th	10Th	11Th	12Th
Hf-175		1.39(36)E-27		
Ta-176	9.6(27)E-28	1.17(12)E-27		
Ta-178m	1.21(17)E-28	1.21(11)E-27		
Re-181		1.13(9)E-27		
Os-182		1.14(17)E-27	8.61(66)E-28	
Re-182	3.23(29)E-27			
Re-183		3.81(20)E-27		
Os-183m		5.35(43)E-28		
Ta-185		3.27(73)E-27		
Os-185		1.21(10)E-27		
Ir-186		1.58(19)E-27	9.8(12)E-28	
Ir-186m			1.50(18)E-27	
Pt-186			3.45(71)E-28	
W-187		1.95(70)E-27		
Ir-188		1.17(27)E-27		
Pt-188		1.07(52)E-27		
Pt-191		1.53(29)E-27	1.30(33)E-27	
Au-191		2.01(27)E-27		
Hg-191m		7.3(25)E-28		
Au-192	1.76(22)E-27		1.6(11)E-27	
Hg-192	7.93(70)E-28	6.98(47)E-28	5.77(76)E-28	2.28(31)E-28
Au-193		3.88(54)E-27		
Hg-193		1.50(42)E-27		
Hg-193m		2.73(39)E-28		
Tl-194m		1.27(33)E-27		
Ir-196m		5.83(89)E-28		
Au-196		5.9(19)E-28		
Tl-196	9.14(83)E-28	1.18(13)E-27	4.96(51)E-28	2.20(32)E-28
Tl-198	1.38(18)E-27	1.46(22)E-27		2.98(65)E-28
Tl-198m		8.6(26)E-28		
Pb-198		6.8(21)E-28		
Pb-199	2.87(64)E-27			
Tl-200	7.5(14)E-28		5.07(83)E-28	
Pb-200	5.15(65)E-28	6.92(46)E-28	3.72(34)E-28	1.74(41)E-28
Pb-201		8.22(33)E-28	4.74(33)E-28	2.11(22)E-28
Pb-202m		6.2(14)E-28	6.22(75)E-28	

Table 2 continued

Residual nuclei	Reaction rate ($\text{atom}^{-1} \cdot \text{deuteron}^{-1}$)			
	9Th	10Th	11Th	12Th
Bi-202		1.07(19)E-27		
Po-202	4.50(57)E-27			
Hg-203		1.30(20)E-27		
Bi-203		6.02(79)E-28		
Po-204	4.21(31)E-27			
Bi-204	9.47(53)E-28	1.47(9)E-27	1.02(6)E-27	4.34(31)E-28
Bi-205		1.30(18)E-27		
Po-205		1.65(42)E-27		
Bi-206		8.8(12)E-28		
Po-207	9.8(15)E-28	1.55(14)E-27	8.8(15)E-28	2.88(60)E-28
At-208			5.2(12)E-28	
At-209	1.21(57)E-27	1.87(5)E-27	1.14(10)E-27	5.10(45)E-28
At-210		1.24(9)E-27	7.74(58)E-28	2.89(23)E-28
Rn-211	4.82(74)E-28	8.74(39)E-28		3.69(41)E-28
Bi-213		3.56(23)E-27		
Ac-224		2.99(38)E-27		8.78(69)E-28
Ac-226	1.82(21)E-27	3.31(7)E-27	2.57(29)E-27	1.01(6)E-27
Th-227		4.08(49)E-27		
Th-231	3.73(42)E-26	7.04(26)E-26	4.85(29)E-26	2.91(21)E-26
Pa-233	7.22(23)E-26	1.82(5)E-25	1.24(4)E-25	7.66(23)E-26

The obtained experimental results were compared with Monte Carlo simulations performed with the FLUKA code. Ratio of the experimental and calculated reaction rates for residual nuclei ^{233}Pa for the sample 9Th is 1.64(5), for 10Th is 1.75(5), for 11Th is 1.70(5), and for 12Th is 1.78(5). For residual nuclei ^{231}Th , the ratio of the experimental and calculated reaction rates for the sample 9Th is 1.68(19), for 10Th is 1.95(7), for 11Th is 2.32(14), and for 12Th is 2.50(18). Ratio of the experimental and calculated cumulative reaction rates for such residual nuclei of fission reactions as ^{87}Kr , ^{88}Kr , ^{91}Sr , ^{92}Sr , ^{92}Y , ^{93}Y , ^{97}Zr , ^{103}Ru , ^{105}Ru , ^{105}Rh , ^{115}Cd , ^{128}Sb , ^{132}Te , ^{131}I , ^{132}I , ^{133}I , ^{135}I , ^{135}Xe , ^{140}Ba , ^{142}La , ^{143}Ce are in the range of 1.4–2.8, for such residual nuclei as ^{96}Nb , ^{99}Mo , ^{113}Ag , ^{122}Sb , ^{124}Sb , ^{130}I , ^{135}Ce , ^{141}Ce are in the range of 3.4–6.6. Ratio of the experimental and calculated total fission reaction rates for the sample 9Th is 2.69(39), for 10Th is 1.29(15), for 11Th is 1.78(19), and for 12Th is 1.87(21). Calculations of fission reaction rates $R(n, f)$ from the experimental data were carried out as follows. Cumulative yields Y of the fission products of ^{232}Th at a neutron energy

of 14 MeV were taken from the TENDL-2011 [40] library. The average values of the fission product R/Y ratios for the following nuclei: ^{85m}Kr , ^{87}Kr , ^{88}Kr , ^{91}Sr , ^{92}Y , ^{92}Sr , ^{93}Y , ^{95}Zr , ^{96}Nb , ^{97}Zr , ^{129}Sb , ^{131}I , ^{132}Te , ^{132}I , ^{132}Cs , ^{133}I , ^{134}I , ^{135}I , ^{135}Xe , ^{140}Ba , ^{141}Ce , ^{142}La , and ^{143}Ce were found to be $5.57(80)\text{E-}26$ for 9Th, $9.8(11)\text{E-}26$ for 10Th, $6.91(73)\text{E-}26$ for 11Th, and $3.50(39)\text{E-}26$ for the sample 12Th. These numbers are the fission reaction rates $R(n, f)$ for ^{232}Th .

6. CONCLUSIONS

The QUINTA setup was irradiated by deuterons with energy 6 GeV. The obtained experimental results were compared with the results of simulations by the FLUKA code. By comparison of experimental and calculated data, agreements were found for residual nuclei ^{233}Pa , ^{231}Th and for several products of fission reactions such as ^{87}Kr , ^{88}Kr , ^{91}Sr , ^{92}Sr , ^{92}Y , ^{93}Y , ^{97}Zr , ^{99}Mo , ^{103}Ru , ^{105}Ru , ^{105}Rh , ^{115}Cd , ^{128}Sb , ^{132}Te , ^{131}I , ^{132}I , ^{133}I , ^{135}I , ^{135}Xe , ^{140}Ba , ^{142}La , ^{143}Ce . For other products of fission reactions the ratio of the experimental and calculated cumulative reaction rates is above 3. We suppose these ratios depend on the yield of isotopes in fission reactions induced by high-energy particles. In addition, for isomers the models are not yet able to predict the ground state/metastable level(s) split.

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